

PHASE TRANSITION TO HYPERONS IN PROTO-NEUTRON STARS

Jérôme Novak (Jerome.Novak@obspm.fr)

Laboratoire Univers et Théories (LUTH)
CNRS / Observatoire de Paris / Université Paris-Diderot

Peres, Oertel & Novak,
Phys. Rev. D **87**, 043006 (2013)

NewCompStar WG3 meeting, November, 19th 2014

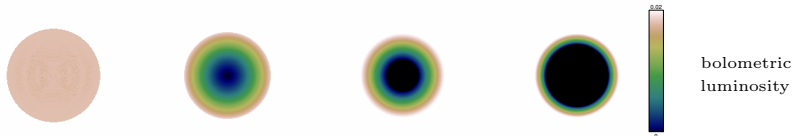
CONTEXT

HOW DO STELLAR BLACK HOLES FORM?

- Stellar-mass black holes form in the collapse of massive stars
- Beginning of collapse triggered by mass-limit of iron core
- Collapse & bounce, then collapse of the proto-neutron star triggered by accretion

⇒ very similar scenario to core-collapse supernova

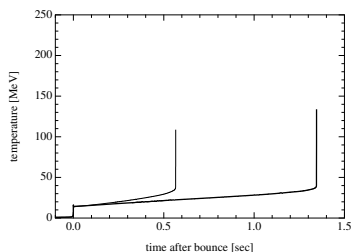
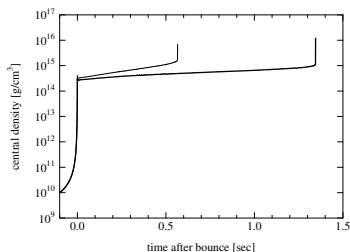
⇒ central engine for gamma-ray bursts (collapsar model)



Vincent *et al.* (2012)

CONTEXT

Collapse to black hole from stellar progenitor has already been studied (*e.g.* Sumiyoshi *et al.* (2007), Fischer *et al.* (2009), O'Connor & Ott (2011), Ugliano *et al.* (2012)...).



40 M_{\odot} progenitor, from Sumiyoshi *et al.* (2007)

⇒ much higher densities (above nuclear saturation density) and temperatures (tens of MeV) than in supernova simulations.

AIMS...

High density & temperature conditions \Rightarrow additional particles should appear (observed on Earth).

- How many “exotic” particles could appear on the way to the black hole?
- What is their influence on the collapse?
- What is their observational signature?

Reverse question:

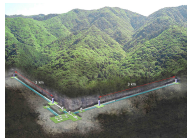
- Can we infer nuclear matter composition from observations of black hole formation?



LIGO



Virgo



KAGRA

PHYSICAL FRAMEWORK

- Spherical or axial symmetry (1D/2D runs).
- Relativistic hydrodynamics, with perfect-fluid stress-energy tensor.
- General relativity in 3+1 formulation. Isotropic gauge for 1D, conformally-flat condition (CFC) in 2D.
- Apparent horizon finder (Lin & Novak 2007).
- Microphysical equation of state from Oertel *et al.* (2012).
- Deleptonization and neutrino leakage.
- Gravitational waves extracted with the modified quadrupole formula (2D).

NUMERICAL TOOLS

CoCoNuT code (Dimmelmeier *et al.* 2005):



- Potentially 3D code, but used only in 1D or 2D (not fully parallel, yet);
- high resolution-shock capturing schemes for the relativistic hydrodynamics (*e.g.* Font 2008) \Rightarrow conservative-form hydrodynamic equations;
- multi-domain pseudo-spectral methods for the solution of Einstein equations (*e.g.* Grandclément & Novak 2009) \Rightarrow non-linear coupled elliptic system;
- interpolation and filtering to avoid Gibbs phenomenon.

EQUATION OF STATE

OERTEL *et al.* (2012), GULMINELLI *et al.* (2012)

Earth-based experiments \Rightarrow pions and hyperons at high densities and temperatures.

Nucleonic interaction from Lattimer & Swesty (1991) \Rightarrow $n, p, e^-, e^+, \gamma, \alpha, A$

EOS LS220+PIONS

- Pions π^-, π^0, π^+
- free gas

EOS LS220+HYPERONS

- Λ hyperons
- contains a first order **phase transition** to hyperonic matter

Hadronic interaction different from previous studies (RMF) with additional particles (Sumiyoshi *et al.* 2009, Shen *et al.* 2011).

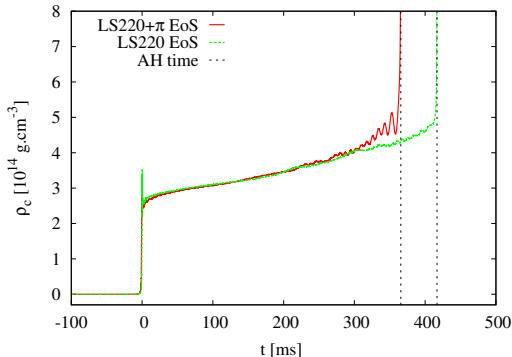
\Rightarrow Compatible with $\sim 2 M_{\odot}$ neutron star observations

Demorest *et al.* (2010), Antoniadis *et al.* (2013).

Results

SPHERICAL SYMMETRY

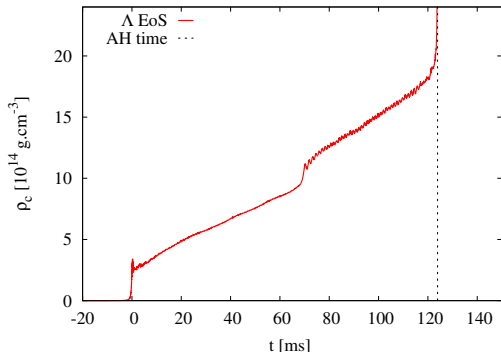
WITH PIONS



- Maximum pion fraction $Y_{\pi^-} = 0.13$ at the onset of BH collapse
- $Y_{\pi^-} > Y_{\pi^0} > Y_{\pi^+}$ at all times

- The PNS with LS220+ π EoS is slightly more compressible
- More compressible means less pressure \rightarrow cannot hold as much mass as LS220 \rightarrow less time post bounce accreting mass and maximum mass smaller
- PNS baryonic masses at BH collapse : $2.55 M_{\odot}$ with LS220, $2.49 M_{\odot}$ with LS220+ π

SPHERICAL SYMMETRY

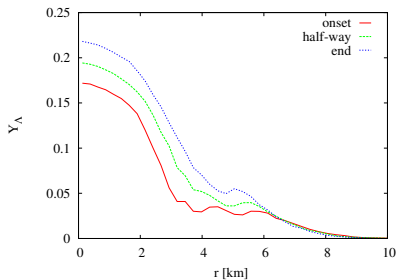
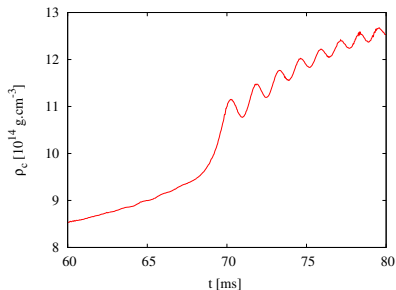


WITH HYPERONS

- PNS baryonic mass at BH collapse : $2.00 M_{\odot}$
- Maximum of Λ hyperon fraction $Y_{\Lambda} = 0.41$ at the onset of BH collapse
- Presence of a **phase transition** to hyperonic matter (related to the high accretion rate)
- The PNS oscillates after the phase transition (PNS fundamental modes)
- Oscillations are resolved in time and stay when increasing the resolution

PHASE TRANSITION

SPHERICAL SYMMETRY

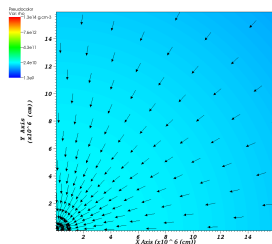
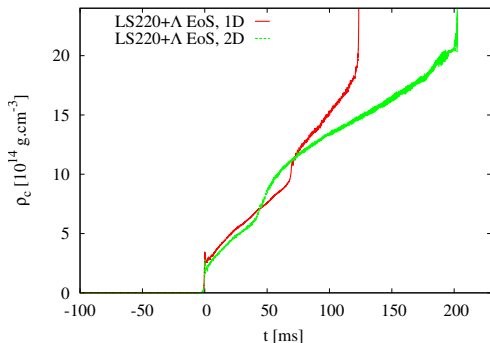


- Phase transition only reached for progenitors with high mass accretion rate (low metallicity),
- Induces a “mini-collapse” followed by oscillations of the PNS,
- No second shock wave as in simulations with phase transition to quark matter (Sagert *et al.* 2009).

PHASE TRANSITION

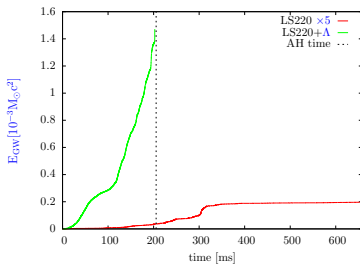
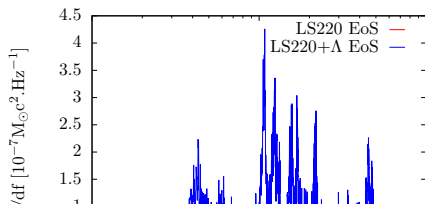
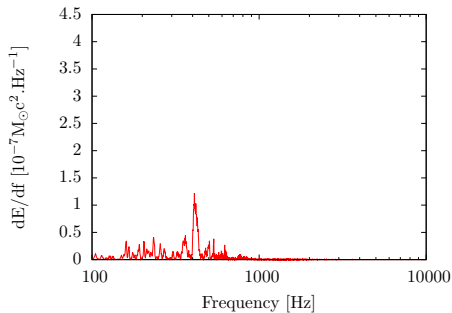
ROTATIONAL SYMMETRY

- 2D in axisymmetry
- Progenitor rotation profile : slow and differential
- All other settings similar to 1D settings



GRAVITATIONAL WAVES

With the modified quadrupole formula



SUMMARY / OUTLOOK

- EoS for core-collapse based on Lattimer & Swesty (1991), with additional particles (π, Λ), compatible with recent observations of $2M_{\odot}$ neutron stars.
 - Softens the PNS, which collapses more rapidly and eventually undergoes a **phase transition** to hyperonic matter.
 - Phase transition “softened” in 2D simulations
⇒ implications for QGP phase transition?
 - Possibly observable with gravitational waves.
-
- Improvement of resolution in 2D
 - Better (full?) neutrino transport (Peres *et al.* 2014)

REFERENCES

-  Antoniadis, J. *et al.*, *Science* **340**, p.448 (2013)
-  <http://www.mpa-garching.mpg.de/hydro/COCGNUT>
-  Demorest, P. *et al.*, *Nature* **467**, p.1081 (2010)
-  Dimmelmeier, H. *et al.*, *Phys. Rev. D* **71**, 064023 (2005)
-  Fischer, T. *et al.*, *Astron. Astrophys.* **499**, 1 (2009)
-  Font, J.A., *Living Rev. Relat.* **11**, 7 (2008)
-  Grandclément, P. and Novak, J., *Living Rev. Relat.* **12**, 1, (2009)
-  Gulminelli, F. *et al.*, *Phys. Rev. C* **86**, 025805 (2012)
-  Lattimer, J.M. and Swesty, F.D., *Nucl. Phys. A* **535**, p.331 (1991)
-  Lin, L.-M. and Novak, J., *Classical Quantum Gravity* **24**, p.2665 (2007)
-  O'Connor, E. and Ott, C.D., *Astrophys. J.* **730**, p.70 (2011)
-  Oertel, M. *et al.*, *Phys. Rev. C* **85**, 055806 (2012)
-  Peres, B. *et al.*, *Phys. Rev. D* **87**, 043006 (2013)
-  Peres, B. *et al.*, *Class. Quantum Grav.* **31** 045012 (2014)
-  Sagert, I. *et al.*, *Phys. Rev. Lett.* **102**, 081101 (2009)
-  Shen, H. *et al.*, *Nucl. Phys. A* **637**, p.435 (1998)
-  Shen, H. *et al.*, *Astrophys. J. Suppl. Ser.* **197**, p.20 (2011)
-  Sumiyoshi *et al.*, *Astrophys. J.* **667**, p.382 (2007)
-  Sumiyoshi *et al.*, *Astrophys. J. Lett.* **690**, L43 (2009)
-  Ugliano *et al.*, *Astrophys. J.* **757**, p.69 (2012)
-  Vincent, F.H. *et al.*, *Classical Quantum Gravity* **29**, 245005 (2012)
-  Woosley *et al.*, *Rev. Mod. Phys.* **74**, p.1015 (2002)

NEUTRINO LEAKAGE

- Only one **opaque** (\Rightarrow fluid) zone and one **transparent** (\Rightarrow free-streaming) zone (*e.g.* van Riper *et al.* 1981)
- No transport, cheap in CPU time, but number of **approximations** and drawbacks
- No semi-transparent regime, no self-consistent heating \Rightarrow not good to revive the shock.

\Rightarrow computation of “optical” depth for three species of neutrinos: $\nu_e, \bar{\nu}_e, \nu_x$. Loss of energy & momentum taken into account.

CREATION PROCESSES

- $p + e^- \rightarrow \nu_e + n$
- $(A, Z) + e^- \rightarrow (A, Z-1) + \nu_e$
- $e^- + e^+ \rightarrow \nu_i + \bar{\nu}_i$
- $\tilde{\gamma} \rightarrow \nu_i + \bar{\nu}_i$

OPACITY PROCESSES

- $\nu_i + N \rightarrow \nu_i + N$
- $\nu_i + (A, Z) \rightarrow \nu_i + (A, Z)$
- $\nu_e + n \rightarrow p + e^-$
- $\bar{\nu}_e + p \rightarrow n + e^+$

Parameters & Initial models

INITIAL SETUP

PROGENITOR

- From Woosley *et al.* (2002), $40M_{\odot}$
ZAMS and $10^{-4} \times$ solar metallicity

LEAKAGE

- β -equilibrium density $1.2 \times 10^{12} \text{ g.cm}^{-2}$
- ν escape time $t_{esc} = 3(R_{\nu\text{-sphere}} - r)\tau$
- power lost by the fluid in the trapped regime $Q_E = -1.1 \langle \epsilon_{\nu} \rangle \frac{Y_{\nu}}{t_{esc}}$

EoS

- Values of the parameters for $Y - N$ and $Y - Y$ interactions compatible with hyperonic data and PSR J 1614-2230 (marginally).